

Volume and Participating Media

Digital Image Synthesis

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with slides by Pat Hanrahan and Torsten Moller

Participating media



- We have by far assumed that the scene is in a vacuum. Hence, radiance is constant along the ray. However, some real-world situations such as fog and smoke attenuate and scatter light. They participate in rendering.

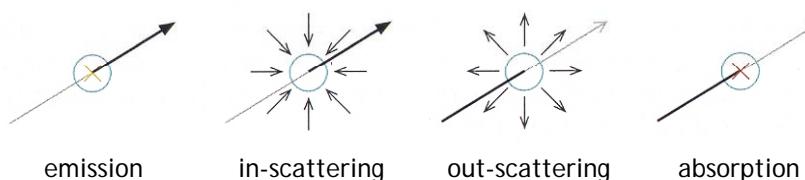
- Natural phenomena
 - Fog, smoke, fire
 - Atmosphere haze
 - Beam of light through clouds
 - Subsurface scattering



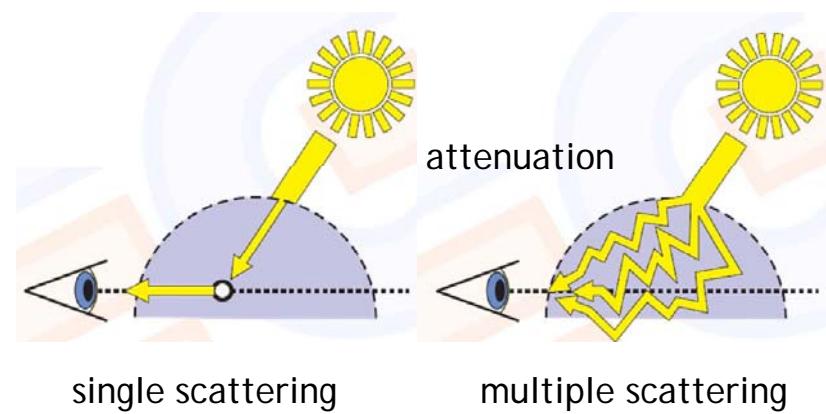
Volume scattering processes



- Absorption (conversion from light to other forms)
- Emission (contribution from luminous particles)
- Scattering (direction change of particles)
 - Out-scattering
 - In-scattering
 - Single scattering v.s. multiple scattering
- Homogeneous v.s. inhomogeneous(heterogeneous)



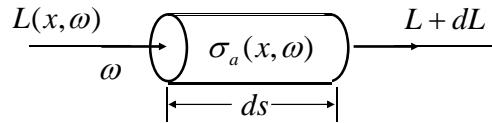
Single scattering and multiple scattering



Absorption



The reduction of energy due to conversion of light to another form of energy (e.g. heat)



$$dL(x, \omega) = -\sigma_a(x, \omega)L(x, \omega)ds$$

Absorption cross-section: $\sigma_a(x, \omega)$

Probability of being absorbed per unit length

Transmittance and opacity



$$dL(x, \omega) = -\sigma_a(x, \omega)L(x, \omega)ds$$

$$\frac{dL(x, \omega)}{L(x, \omega)} = -\sigma_a(x, \omega)ds \rightarrow \int_x^{x+s\omega} \frac{dL(x', \omega)}{L(x', \omega)} = -\int_0^s \sigma_a(x + s'\omega, \omega)ds'$$

$$\ln L(x + s\omega, \omega) - \ln L(x, \omega) = -\int_0^s \sigma_a(x + s'\omega, \omega)ds' = -\tau_\omega(s)$$

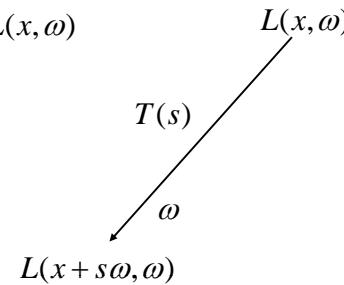
$$L(x + s\omega, \omega) = e^{-\tau_\omega(s)}L(x, \omega) = T_\omega(s)L(x, \omega)$$

Transmittance

$$T_\omega(s) = e^{-\tau_\omega(s)}$$

Opacity

$$\alpha_\omega(s) = 1 - T_\omega(s)$$



Transmittance



$$dL(x, \omega) = -\sigma_a(x, \omega)L(x, \omega)ds$$

$$\frac{dL(x, \omega)}{L(x, \omega)} = -\sigma_a(x, \omega)ds \rightarrow \int_x^{x+s\omega} \frac{dL(x', \omega)}{L(x', \omega)} = -\int_0^s \sigma_a(x + s'\omega, \omega)ds'$$

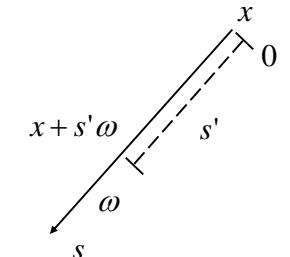
$$\ln L(x + s\omega, \omega) - \ln L(x, \omega) = -\int_0^s \sigma_a(x + s'\omega, \omega)ds' = -\tau_\omega(s)$$

Optical distance or depth

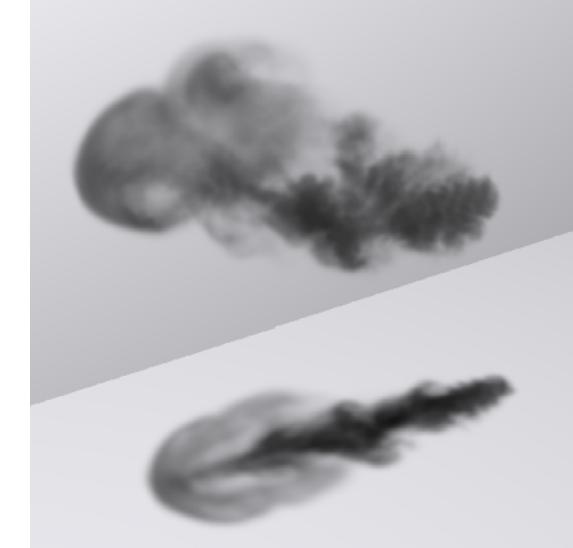
$$\tau_\omega(s) = \int_0^s \sigma_a(x + s'\omega, \omega)ds'$$

Homogenous media: constant σ_a

$$\sigma_a \rightarrow \tau(s) = \sigma_a s$$



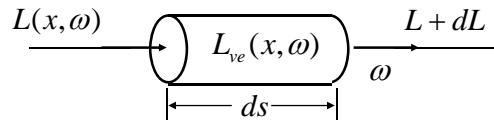
Absorption



Emission

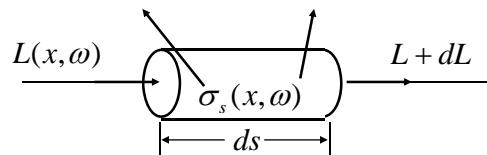
- Energy that is added to the environment from luminous particles due to chemical, thermal, or nuclear processes that convert energy to visible light.
- $L_{ve}(x, \omega)$: emitted radiance added to a ray per unit distance at a point x in direction ω

$$dL(x, \omega) = L_{ve}(x, \omega)ds$$



Out-scattering

Light heading in one direction is scattered to other directions due to collisions with particles

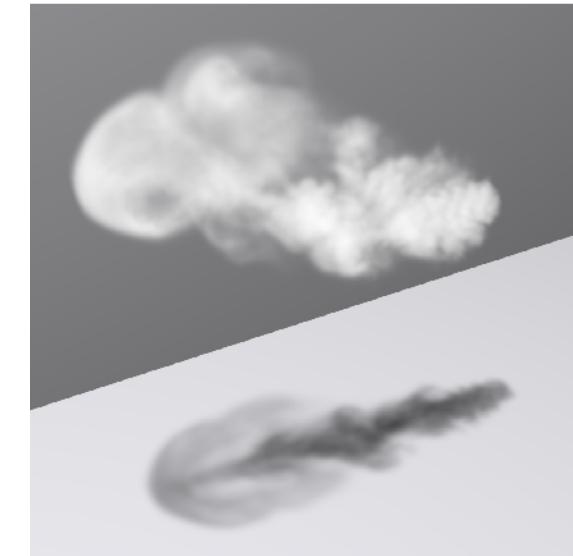


$$dL(x, \omega) = -\sigma_s(x, \omega)L(x, \omega)ds$$

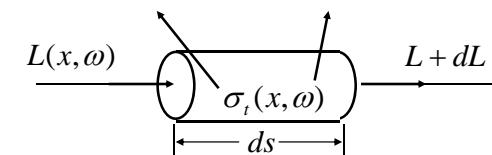
Scattering cross-section: σ_s

Probability of being scattered per unit length

Emission



Extinction



$$dL(x, \omega) = -\sigma_t(x, \omega)L(x, \omega)ds$$

Total cross-section

$$\sigma_t = \sigma_a + \sigma_s$$

Albedo

$$W = \frac{\sigma_s}{\sigma_t} = \frac{\sigma_s}{\sigma_a + \sigma_s}$$

Attenuation due to both absorption and scattering

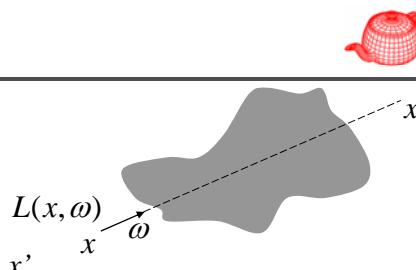
$$\tau_\omega(s) = \int_0^s \sigma_t(x + s' \omega, \omega) ds'$$

Extinction

- Beam transmittance

$$Tr(x \rightarrow x') = e^{-\int_0^s \sigma_t(x+s'\omega, \omega) ds'}$$

s: distance between x and x'



- Properties of Tr :

- In vacuum $Tr(x \rightarrow x') = 1$

- Multiplicative $Tr(x \rightarrow x'') = Tr(x \rightarrow x') \cdot Tr(x' \rightarrow x'')$

- Beer's law (in homogeneous medium)

$$Tr(x \rightarrow x') = e^{-\sigma_t s}$$

Source term



$$S(x, \omega) = L_{ve}(x, \omega) + \sigma_s(x, \omega) \int_{\Omega} p(x, \omega' \rightarrow \omega) L(x, \omega') d\omega'$$

$$dL(x, \omega) = S(x, \omega) ds$$

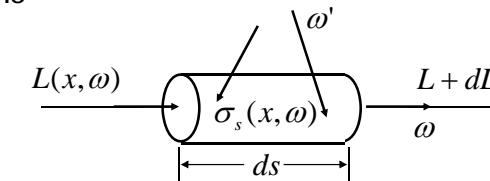
- S is determined by

- Volume emission
- Phase function which describes the angular distribution of scattered radiation (volume analog of BSDF for surfaces)

In-scattering



Increased radiance due to scattering from other directions



$$dL(x, \omega) = \left[\sigma_s(x, \omega) \int_{\Omega} p(x, \omega' \rightarrow \omega) L(x, \omega') d\omega' \right] ds$$

Phase function $p(\omega' \rightarrow \omega)$

Reciprocity

$$p(\omega \rightarrow \omega') = p(\omega' \rightarrow \omega)$$

Energy conserving

$$\int_{S^2} p(\omega' \rightarrow \omega) d\omega' = 1$$

Source term



$$S(x, \omega) = L_{ve}(x, \omega) + \sigma_s(x, \omega) \int_{\Omega} p(x, \omega' \rightarrow \omega) L(x, \omega') d\omega'$$

$$dL(x, \omega) = S(x, \omega) ds$$

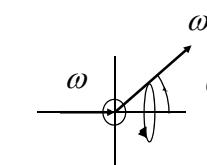
- S is determined by

- Volume emission
- Phase function which describes the angular distribution of scattered radiation (volume analog of BSDF for surfaces)

Phase functions



Phase angle $\cos \theta = \omega \bullet \omega'$



Phase functions

(from the phase of the moon)

1. Isotropic

- simple

$$p(\cos \theta) = \frac{1}{4\pi}$$

2. Rayleigh

- Molecules (useful for very small particles whose radii smaller than wavelength of light)

$$p(\cos \theta) = \frac{3}{4} \frac{1 + \cos^2 \theta}{\lambda^4}$$

3. Mie scattering

- small spheres (based on Maxwell's equations; good model for scattering in the atmosphere due to water droplets and fog)

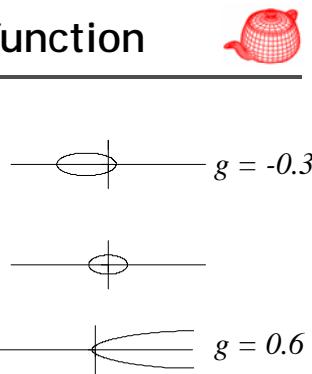
Henyey-Greenstein phase function

Empirical phase function

$$p(\cos \theta) = \frac{1}{4\pi} \frac{1-g^2}{(1+g^2 - 2g \cos \theta)^{3/2}}$$

$$2\pi \int_0^\pi p(\cos \theta) \cos \theta d\theta = g$$

g: average phase angle



Henyey-Greenstein approximation

- Any phase function can be written in terms of a series of Legendre polynomials (typically, n<4)

$$p(\cos \theta) = \frac{1}{4\pi} \sum_{n=0}^{\infty} (2n+1) b_n P_n(\cos \theta)$$

$$\begin{aligned} b_n &= \langle p(\cos \theta), P_n(\cos \theta) \rangle \\ &= \int_{-1}^1 p(\cos \theta) P_n(\cos \theta) d\cos \theta \end{aligned}$$

$$P_0(x) = 1$$

$$P_1(x) = x$$

$$P_2(x) = \frac{1}{2}(3x^2 - 1)$$

$$P_3(x) = \frac{1}{2}(5x^3 - 3x)$$

...

Schlick approximation

- Approximation to Henyey-Greenstein

$$p_{Schlick}(\cos \theta) = \frac{1}{4\pi} \frac{1-k^2}{(1-k \cos \theta)^2}$$

- K plays a similar role like g
 - 0: isotropic
 - 1: back scattering
 - Could use $k = 1.55g - 0.55g^2$



Importance sampling for HG

$$p(\cos \theta) = \frac{1}{4\pi} \frac{1-g^2}{(1+g^2 - 2g \cos \theta)^{3/2}}$$

$$\phi = 2\pi\xi$$

$$\cos \theta = \begin{cases} 1-2\xi & \text{if } g=0 \\ -\frac{1}{|2g|} \left(1+g^2 - \left(\frac{1-g^2}{1-g+2g\xi} \right)^2 \right) & \text{otherwise} \end{cases}$$



Pbrt implementation

```
• core/volume.* volume/*  
class VolumeRegion {  
public:  
    bool IntersectP(Ray &ray, float *t0, float *t1);  
    Spectrum sigma_a(Point &, Vector &);  
    Spectrum sigma_s(Point &, Vector &);  
    Spectrum Lve(Point &, Vector &);  
    // phase functions: pbrt has isotropic, Rayleigh,  
    // Mie, HG, Schlick  
    virtual float p(Point &, Vector &, Vector &);  
    // attenuation coefficient; s_a+s_s  
    Spectrum sigma_t(Point &, Vector &);  
    // calculate optical thickness by Monte Carlo or  
    // closed-form solution  
    Spectrum Tau(Ray &ray, float step=1.,  
                 float offset=0.5);  
};  
    t0   step   t1  
    ●-----●-----●  
    offset
```



Homogenous volume



Homogenous volume

- Determined by (constant)
 - σ_s and σ_a
 - g in phase function
 - Emission L_{ve}
 - Spatial extent

```
Spectrum Tau(Ray &ray, float, float){  
    float t0, t1;  
    if (!IntersectP(ray,&t0,&t1))  
        return 0.;  
    return Distance(ray(t0),ray(t1)) * (sig_a + sig_s);  
}
```



Varying-density volumes

- Density is varying in the medium and the volume scattering properties at a point is the product of the density at that point and some baseline value.
- **DensityRegion**
 - 3D grid, **VolumeGrid**
 - Exponential density, **ExponentialDensity**

DensityRegion

```
class DensityRegion : public VolumeRegion {  
public:  
    DensityRegion(Spectrum &sig_a, Spectrum &sig_s,  
        float g, Spectrum &Le, Transform &VolumeToWorld);  
    float Density(Point &Pobj) const = 0;  
    Spectrum sigma_a(Point &p, Vector &);  
    return Density(WorldToVolume(p)) * sig_a; }  
    Spectrum sigma_s(Point &p, Vector &);  
    return Density(WorldToVolume(p)) * sig_s; }  
    Spectrum sigma_t(Point &p, Vector &);  
    return Density(WorldToVolume(p))*(sig_a+sig_s); }  
    Spectrum Lve(Point &p, Vector &);  
    return Density(WorldToVolume(p)) * le; }  
...  
protected:  
    Transform WorldToVolume;  
    Spectrum sig_a, sig_s, le;  
    float g;  
};
```



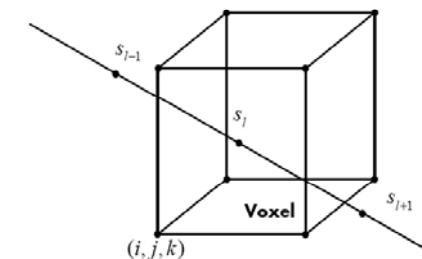
VolumeGrid

```
VolumeGrid(Spectrum &sa, Spectrum &ss, float gg,  
    Spectrum &emit, BBox &e, Transform &v2w,  
    int nx, int ny, int nz, const float *d);  
  
float VolumeGrid::Density(const Point &Pobj) const {  
    if (!extent.Inside(Pobj)) return 0;  
    // Compute voxel coordinates and offsets  
    float vox_x = (Pobj.x - extent.pMin.x) /  
        (extent.pMax.x - extent.pMin.x) * nx - .5f;  
    float vox_y = (Pobj.y - extent.pMin.y) /  
        (extent.pMax.y - extent.pMin.y) * ny - .5f;  
    float vox_z = (Pobj.z - extent.pMin.z) /  
        (extent.pMax.z - extent.pMin.z) * nz - .5f;
```

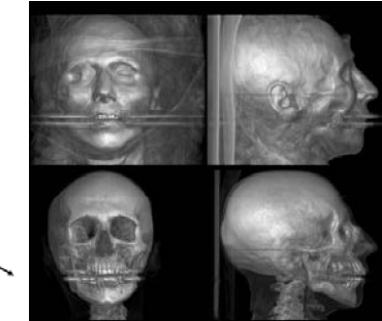


VolumeGrid

- Standard form of given data
- Tri-linear interpolation of data to give continuous volume
- Often used in volume rendering



Interpolation $v(s_j) = \text{trilinear}(v, i, j, k, x(s_j))$



VolumeGrid

```
int vx = Floor2Int(vox_x);  
int vy = Floor2Int(vox_y);  
int vz = Floor2Int(vox_z);  
float dx = vox_x - vx, dy = vox_y - vy, dz = vox_z - vz;  
// Trilinearly interpolate density values  
float d00 = Lerp(dx, D(vx, vy, vz), D(vx+1, vy, vz));  
float d10 = Lerp(dx, D(vx, vy+1, vz), D(vx+1, vy+1, vz));  
float d01 = Lerp(dx, D(vx, vy, vz+1), D(vx+1, vy, vz+1));  
float d11 = Lerp(dx, D(vx, vy+1, vz+1), D(vx+1, vy+1, vz+1));  
float d0 = Lerp(dy, d00, d10);  
float d1 = Lerp(dy, d01, d11);  
return Lerp(dz, d0, d1);  
}  
float D(int x, int y, int z) {  
    x = Clamp(x, 0, nx-1);  
    y = Clamp(y, 0, ny-1);  
    z = Clamp(z, 0, nz-1);  
    return density[z*nx*ny+y*nx+x];  
}
```



Exponential density

- Given by

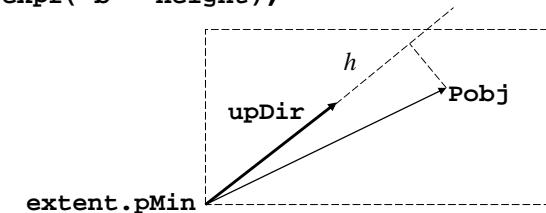
$$d(h) = ae^{-bh}$$

- Where h is the height in the direction of the up-vector



ExponentialDensity

```
class ExponentialDensity : public DensityRegion {
public:
    ExponentialDensity(Spectrum &sa, Spectrum &ss,
                      float g, Spectrum &emit, BBox &e, Transform &v2w,
                      float aa, float bb, Vector &up)
    ...
    float Density(const Point &Pobj) const {
        if (!extent.Inside(Pobj)) return 0;
        float height = Dot(Pobj - extent.pMin, upDir);
        return a * expf(-b * height);
    }
private:
    BBox extent;
    float a, b;
    Vector upDir;
};
```



Light transport



- Emission + in-scattering (source term)

$$S(x, \omega) = L_{ve}(x, \omega) + \sigma_s(x, \omega) \int_{\Omega} p(x, \omega' \rightarrow \omega) L(x, \omega') d\omega'$$

$$dL(x, \omega) = S(x, \omega) ds$$

- Absorption + out-scattering (extinction)

$$dL(x, \omega) = -\sigma_t(x, \omega) L(x, \omega) ds$$

- Combined

$$\frac{dL(x, \omega)}{ds} = -\sigma_t(x, \omega) L(x, \omega) + S(x, \omega)$$

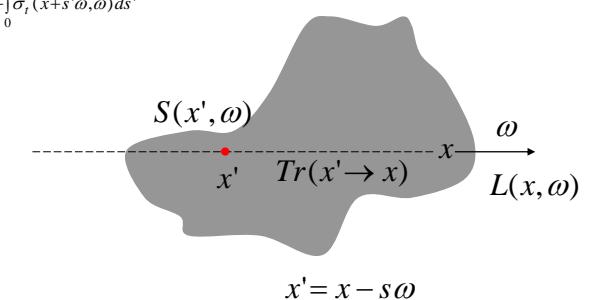
Infinite length, no surface



- Assume that there is no surface and we have an infinite length, we have the solution

$$L(x, \omega) = \int_0^{\infty} Tr(x' \rightarrow x) S(x', \omega) ds'$$

$$Tr(x' \rightarrow x) = e^{-\int_0^s \sigma_t(x+s', \omega) ds'}$$

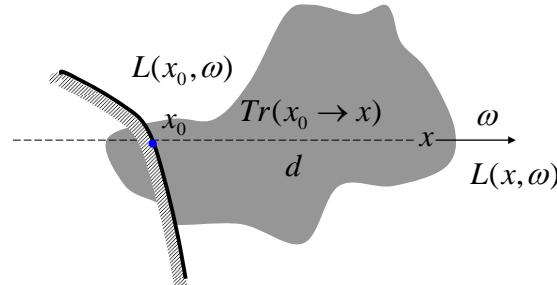


With surface

- The solution

$$L(x, \omega) = Tr(x_0 \rightarrow x)L(x_0, -\omega)$$

from the surface point x_0

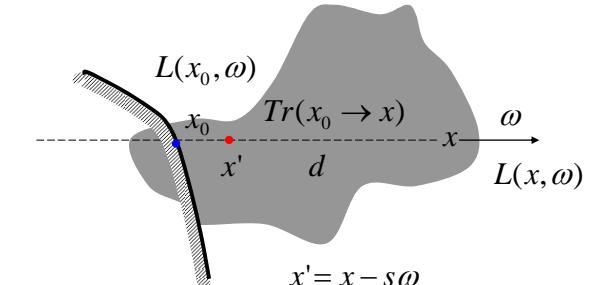


With surface

- The solution

$$L(x, \omega) = Tr(x_0 \rightarrow x)L(x_0, -\omega) + \int_0^d Tr(x' \rightarrow x)S(x', -\omega)ds$$

from the surface point x_0 from the participating media



$$x' = x - s \omega$$

Simple atmosphere model

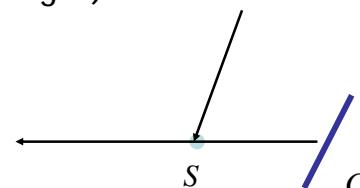


Assumptions

- Homogenous media
- Constant source term (airlight)

$$\frac{\partial L(s)}{\partial s} = -\sigma_i L(s) + S$$

$$L(s) = (1 - e^{-\sigma_i s})S + e^{-\sigma_i s}C$$

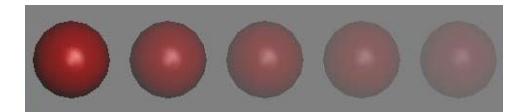


- Fog
- Haze

OpenGL fog model



$$C = fC_{in} + (1 - f)C_{fog}$$



GL_EXP

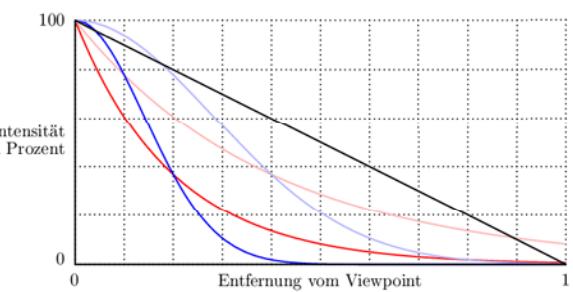
$$f(z) = e^{-(density \cdot z)}$$

GL_EXP2

$$f(z) = e^{-(density \cdot z)^2}$$

GL_LINEAR

$$f(z) = \frac{end - z}{end - start}$$



From <http://wiki.delphigl.com/index.php/glFog>

VolumeIntegrator



```
class VolumeIntegrator : public Integrator {  
public:  
    Beam transmittance for a given  
    ray from mint to maxt  
    virtual Spectrum Transmittance(  
        const Scene *scene,  
        const Renderer *renderer,  
        const RayDifferential &ray,  
        const Sample *sample, ...) const = 0;  
};
```

Pick up functions `Preprocess()`, `RequestSamples()` and `Li()` from `Integrator`.

Emission only



- Use multiplicativity of Tr

$$Tr(x_i \rightarrow x) = Tr(x_i \rightarrow x_{i-1}) \cdot Tr(x_{i-1} \rightarrow x)$$

- Break up integral and compute it incrementally by ray marching
 - Tr can get small in a long ray
 - Early ray termination
 - Either use Russian Roulette or deterministically

Emission only



- Solution for the emission-only simplification

$$S(x', -\omega) = L_{ev}(x', -\omega)$$

$$L(x, \omega) = Tr(x_0 \rightarrow x)L(x_0, -\omega) + \int_0^d Tr(x' \rightarrow x)L_{ev}(x', -\omega)ds$$

- Monte Carlo estimator

$$\frac{1}{N} \sum_{i=1}^N \frac{\text{Tr}(x_i \rightarrow x) L_{ev}(x_i, \omega)}{p(x_i)} = \frac{t_1 - t_0}{N} \sum_{i=1}^N \text{Tr}(x_i \rightarrow x) L_{ev}(x_i, \omega)$$

EmissionIntegrator



```
class EmissionIntegrator : public VolumeIntegrator {
public:
    EmissionIntegrator(float ss) { stepSize = ss; }
    void RequestSamples(Sampler *sampler, Sample
        *sample, Scene *scene);
    Spectrum Li(Scene *scene, Renderer *renderer,
        RayDifferential &ray, Sample *sample, RNG &rng,
        Spectrum *transmittance, MemoryArena &arena);
    Spectrum Transmittance(Scene *scene, Renderer *,
        RayDifferential &ray, Sample *sample, RNG &rng,
        MemoryArena &arena);
private:
    float stepSize;
    int tauSampleOffset, scatterSampleOffset;
};
```

EmissionIntegrator::Transmittance

```

if (!scene->volumeRegion) return Spectrum(1);
float step, offset;
if (sample) { smple is non-NULL only for camera rays
    step = stepSize;
    offset = sample->oneD[tauSampleOffset][0];
} else {
    step = 4.f * stepSize; use larger steps for shadow and
    indirect rays for efficiency
    offset = rng.RandomFloat();
}
Spectrum tau = scene->volumeRegion->tau(ray,
                                             step, offset);
return Exp(-tau);           
$$\tau_{\omega}(s) = \int_0^s \sigma_a(x + s' \omega, \omega) ds'$$


$$T_{\omega}(s) = e^{-\tau_{\omega}(s)}$$


```



EmissionIntegrator::Li

```

VolumeRegion *vr = scene->volumeRegion;
float t0, t1;
if (!vr || !vr->IntersectP(ray, &t0, &t1)
    || (t1-t0) == 0.f) {
    *T = Spectrum(1.f);
    return 0.f;
}
// Do emission-only volume integration in vr
Spectrum Lv(0.);
// Prepare for volume integration stepping
int nSamples = Ceil2Int((t1-t0) / stepSize);
float step = (t1 - t0) / nSamples;
Spectrum Tr(1.f);
Point p = ray(t0), pPrev;
Vector w = -ray.d;
t0 += sample->oneD[scatterSampleOffset][0] * step;

```



EmissionIntegrator::Li

```

for (int i = 0; i < nSamples; ++i, t0 += step) {
    pPrev = p; p = ray(t0);
    Ray tauRay(pPrev, p - pPrev, 0.f, 1.f, ray.time,
               ray.depth);
    Spectrum stepTau = vr->tau(tauRay,.5f * stepSize,
                                  rng.RandomFloat());
    Tr *= Exp(-stepTau); 
$$Tr(x_i \rightarrow x) = Tr(x_i \rightarrow x_{i-1}) \cdot Tr(x_{i-1} \rightarrow x)$$

    // Possibly terminate if transmittance is small
    if (Tr.y() < 1e-3) {
        const float continueProb = .5f;
        if (rng.RandomFloat() > continueProb) break;
        Tr /= continueProb;
    }
    // Compute emission-only source term at _p_
    Lv += Tr * vr->Lve(p, w, ray.time);
}
*T = Tr;           
$$T = \frac{t_1 - t_0}{N} \sum_{i=1}^N Tr(x_i \rightarrow x) L_{ev}(x_i, \omega)$$

return Lv * step;

```



Emission only



exponential density

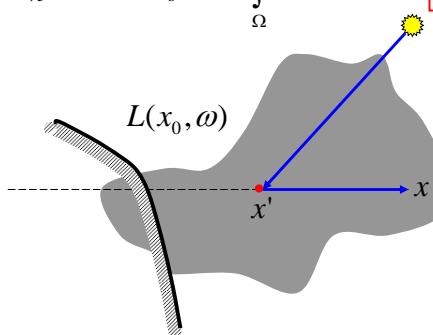


Single scattering

- Consider incidence radiance due to direct illumination

$$L(x, \omega) = Tr(x_0 \rightarrow x)L(x_0, \omega) + \int_0^d Tr(x' \rightarrow x)S(x', \omega)ds$$

$$S(x, \omega) = L_{ve}(x, \omega) + \sigma_s(x, \omega) \int_{\Omega} p(x, \omega' \rightarrow \omega) L_d(x, \omega') d\omega'$$

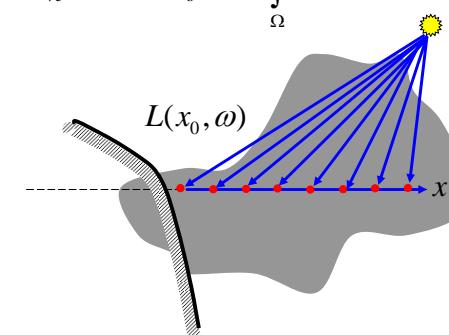


Single scattering

- Consider incidence radiance due to direct illumination

$$L(x, \omega) = Tr(x_0 \rightarrow x)L(x_0, \omega) + \int_0^d Tr(x' \rightarrow x)S(x', \omega)ds$$

$$S(x, \omega) = L_{ve}(x, \omega) + \sigma_s(x, \omega) \int_{\Omega} p(x, \omega' \rightarrow \omega) L_d(x, \omega') d\omega'$$



Single scattering

- L_d may be attenuated by participating media
- At each point of the integral, we could use multiple importance sampling to get

$$\sigma_s(x, \omega) \int_{\Omega} p(x, \omega' \rightarrow \omega) L_d(x, \omega') d\omega'$$

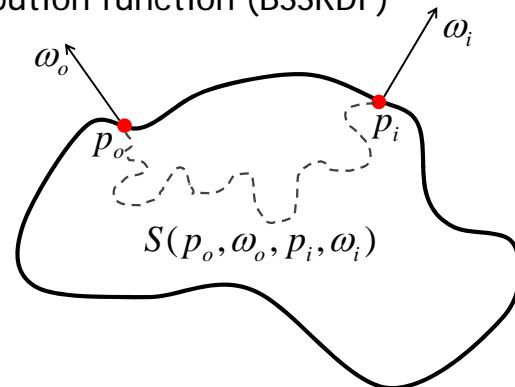
But, in practice, we can just pick up light source randomly.

Single scattering



Subsurface scattering

- The bidirectional scattering-surface reflectance distribution function (BSSRDF)



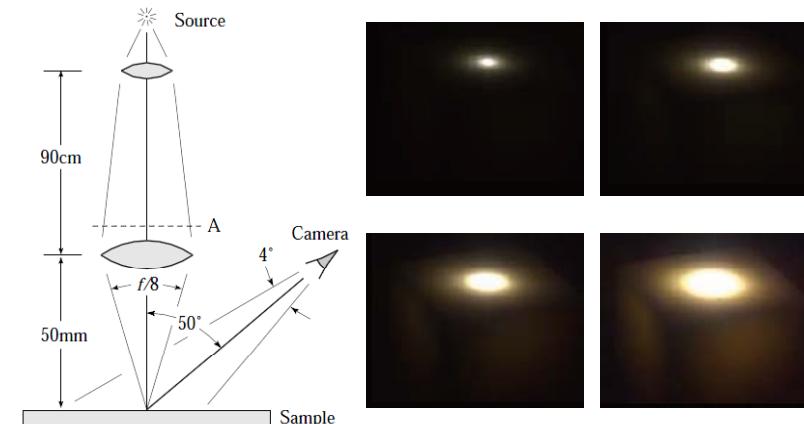
$$L_o(p_o, \omega_o) = \int \int_{A H^2} S(p_o, \omega_o, p_i, \omega_i) L_i(p_i, \omega_i) |\cos \theta_i| d\omega_i dA$$



Subsurface scattering

- Translucent materials have similar mechanism for light scattering as participating media. Thus, path tracing could be used.
- However, many translucent objects have very high albedo. Taken milk as an example, after 100 scattering events, 87.5% of the incident light is still carried by a path, 51% after 500 and 26% after 1,000.
- Efficiently rendering these kinds of translucent scattering media requires a different approach.

Highly translucent materials



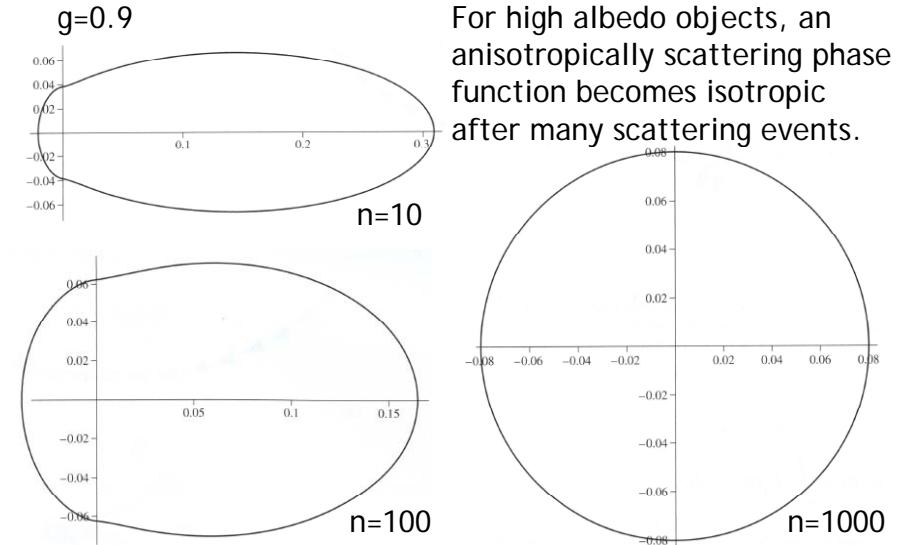
Main idea

- Assume that the material is homogeneous and the medium is semi-infinite, we can use diffusion approximation to describe the equilibrium distribution of illumination.
- There is a solution to the diffusion equation by using a dipole of two light sources to approximate the overall scattering.

Highly scattering media



Principle of similarity



For high albedo objects, an anisotropically scattering phase function becomes isotropic after many scattering events.

Diffusion approximation



- The reduced scattering coefficient

$$\sigma_s' = (1-g)\sigma_s$$

- The reduced extinction coefficient

$$\sigma_t' = \sigma_a + \sigma_s'$$

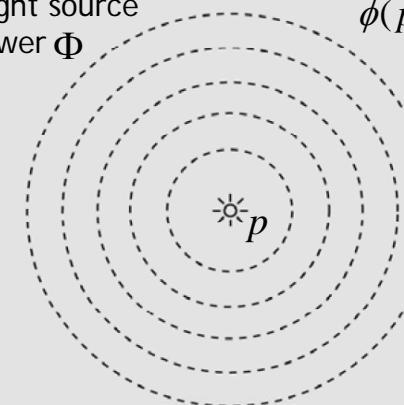
Diffusion approximation



For a point light source
at p with power Φ

$$\phi(p) = \int_{S^2} L_i(p, \omega) d\omega$$

fluence



$$\sigma_{tr} = \sqrt{3\sigma_t' \sigma_a}$$

$$D = 1/(3\sigma_t')$$

diffusion coefficient

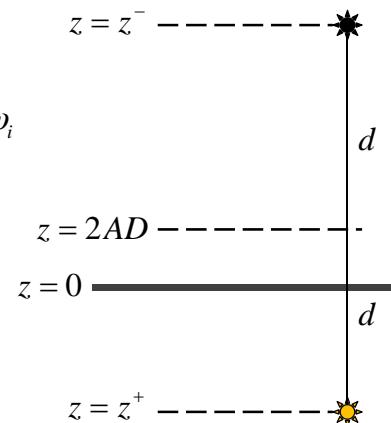
The fluence at p_i from p is

$$\phi(p, p_i) = \frac{\Phi}{4\pi D} \frac{e^{-\sigma_{tr} \|p - p_i\|}}{\|p - p_i\|}$$

Dipole diffusion approximation

$$A = \frac{1 + F_{dr}(\eta)}{1 - F_{dr}(\eta)}$$

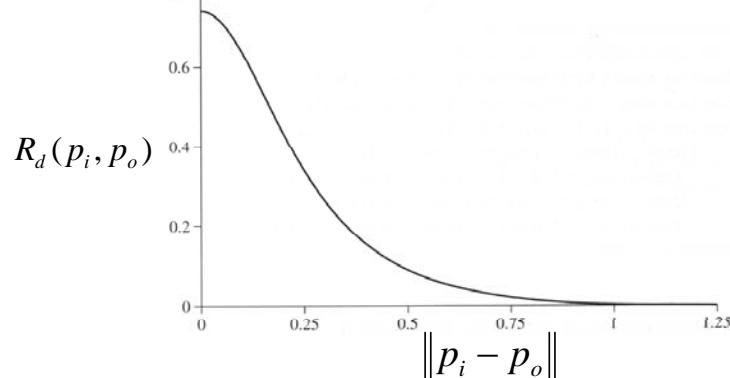
$$F_{dr}(\eta) = \int_{H^2} F_r(\eta, \omega_i) |\omega_i \cdot n| d\omega_i$$



Dipole diffusion approximation

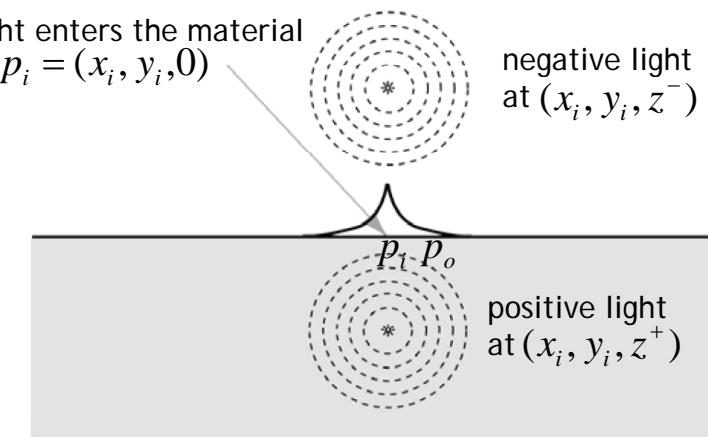
- Put all together

$$R_d(p_i, p_o) = \frac{1}{4\pi} \left(\frac{z^+(d^+ \sigma_{tr} + 1) e^{-\sigma_{tr} d^+}}{(d^+)^3} - \frac{z^-(d^- \sigma_{tr} + 1) e^{-\sigma_{tr} d^-}}{(d^-)^3} \right)$$



Dipole diffusion approximation

Light enters the material at $p_i = (x_i, y_i, 0)$

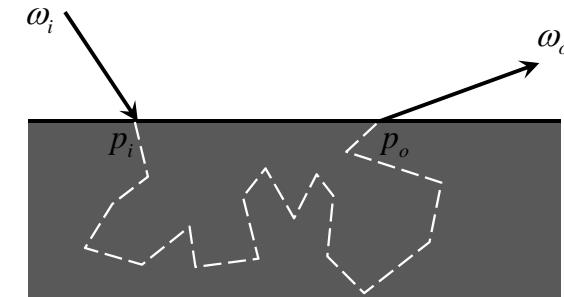


$$\phi(p_i, p_o) = \frac{\Phi}{4\pi D} \left(\frac{e^{-\sigma_{tr} d^+}}{d^+} - \frac{e^{-\sigma_{tr} d^-}}{d^-} \right) \quad d^+ = \| (x_i, y_i, z^+) - p_o \| \quad d^- = \| (x_i, y_i, z^-) - p_o \|$$

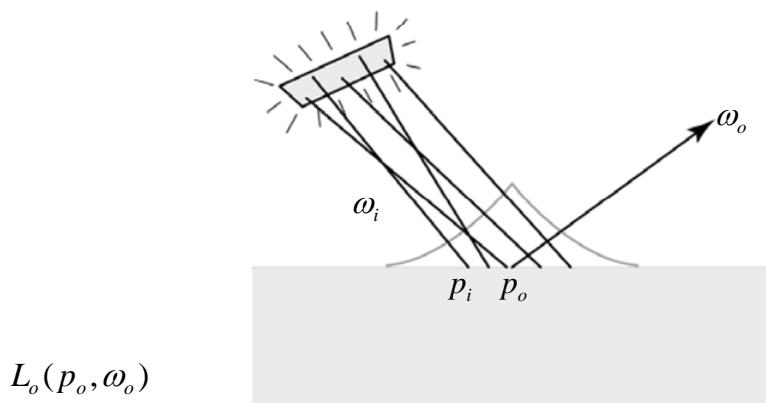
BSSRDF

- BSSRDF based on the diffusion subsurface reflectance approximation

$$S(p_o, \omega_o, p_i, \omega_i) = \frac{1}{\pi} F_t(\eta_o, \omega_o) R_d(\|p_i - p_o\|) F_t(\eta_i, \omega_i)$$

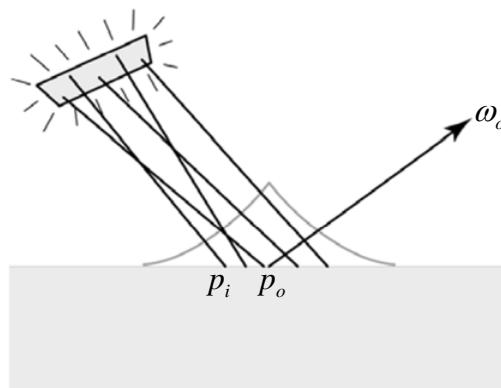


Evaluating BSSRDF



$$L_o(p_o, \omega_o) = \int_A \int_{H^2} \left(\frac{1}{\pi} F_t(\eta_o, \omega_o) R_d(\|p_i - p_o\|) F_t(\eta_i, \omega_i) \right) L_i(p_i, \omega_i) |\cos \theta_i| d\omega_i dA$$

Evaluating BSSRDF



$$L_o(p_o, \omega_o) \approx \frac{1}{\pi} F_t(\eta, \omega_o) F_{dt}(\eta) \int_A R_d(\|p_i - p_o\|) E(p_i) dA$$

$$M_o(p_o) = \int_A R_d(\|p_i - p_o\|) E(p_i) dA \approx \sum_j R_d(\|p_j - p_o\|) E_j A_j$$

Evaluating BSSRDF

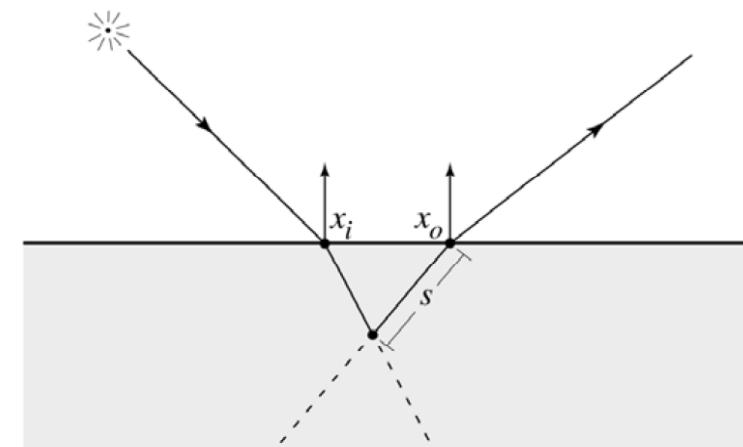


$$L_o(p_o, \omega_o) = \frac{1}{\pi} F_t(\eta_o, \omega_o) \int_A R_d(\|p_i - p_o\|) \left[\int_{H^2} (F_t(\eta_i, \omega_i)) L_i(p_i, \omega_i) |\cos \theta_i| d\omega_i \right] dA \approx F_{dt}(\eta_i) E(p_i)$$

For homogeneous materials, $\eta = \eta_o = \eta_i$

$$L_o(p_o, \omega_o) \approx \frac{1}{\pi} F_t(\eta, \omega_o) F_{dt}(\eta) \int_A R_d(\|p_i - p_o\|) E(p_i) dA$$

Single scattering



Solutions

- Path tracing
- Photon mapping
- Hierarchical approach (Jensen 2002)



Three main components

- Sample a large number of random points on the surface and their incident irradiance is computed.
- Create a hierarchy of these points, progressively clustering nearby points and summing their irradiance values.
- At rendering, use a hierarchical integration algorithm to evaluate the subsurface scattering equation.



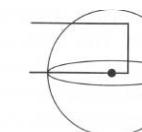
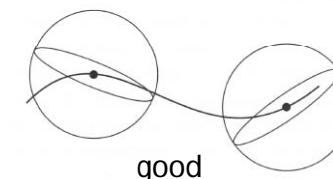
DipoleSubsurfaceIntegrator

```
class DipoleSubsurfaceIntegrator : public
    SurfaceIntegrator {
    ...
    int maxSpecularDepth;
    float maxError, minSampleDist;
    string filename;
    vector<IrradiancePoint> irradiancePoints;
    BBox octreeBounds;
    SubsurfaceOctreeNode *octree;
}
```



Sample points

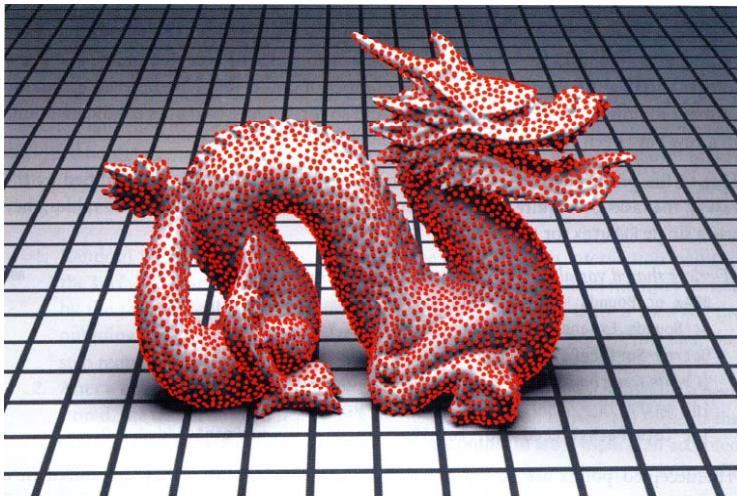
- Samples should be reasonably uniformly distributed. A Poisson disk distribution of points is a good choice.
- There are some algorithms for generating such distributions. Pbrt uses a 3D version of dart throwing by performing Poisson sphere tests. The algorithm terminates when a few thousand tests have been rejected in a row.



Bad, but dipole
is dad for this
anyway



Sample points

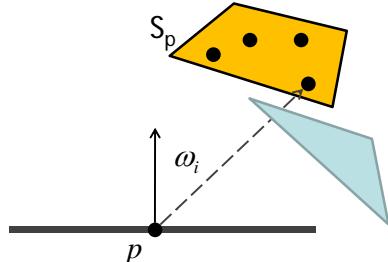


<Compute irradiance values at sample points>

```
for (uint32_t i = 0; i < pts.size(); ++i) {
    SurfacePoint &sp = pts[i];
    Spectrum E(0.f);
    for (int j = 0; j < scene->lights.size(); ++j) {
        <Add irradiance from light at point>
        Calculate direct lighting only; it is possible to include indirect
        lighting but more expensive
        irradiancePoints.push_back(IrradiancePoint(sp, E));
    }
}
```

$$E(p) = \int_{H^2} L_i(p, \omega_i) |\cos \theta_i| d\omega_i$$

$$E(p) = \frac{1}{N} \sum_j \frac{L_i(p, \omega_j) |\cos \theta_j|}{p(\omega_j)}$$



Preprocess

```
if (scene->lights.size() == 0) return;
<Get SurfacePoints for translucent objects>
<Compute irradiance values at sample points>
<Create octree of clustered irradiance samples>
```



IrradiancePoint

```
struct SurfacePoint {
    ...
    Point p;
    Normal n;
    float area, rayEpsilon;
};

struct IrradiancePoint {
    ...
    Point p;
    Normal n;
    Spectrum E;
    float area, rayEpsilon;
};
```



<Create octree of clustered irradiance samples>

```
octree = octreeArena.Alloc<SubsurfaceOctreeNode>();
for (int i = 0; i < irradiancePoints.size(); ++i)
    octreeBounds = Union(octreeBounds,
        irradiancePoints[i].p);
for (int i = 0; i < irradiancePoints.size(); ++i)
    octree->Insert(octreeBounds, &irradiancePoints[i],
        octreeArena);
octree->InitHierarchy();
```

Computes representative position, irradiance and area for each node.
Positions are weighted by irradiance values to emphasize the points
with higher irradiance values.

Li

```
L += UniformSampleAllLights(...);
if (ray.depth < maxSpecularDepth) {
    L += SpecularReflect(...);
    L += SpecularTransmit(...);
}
return L;
```

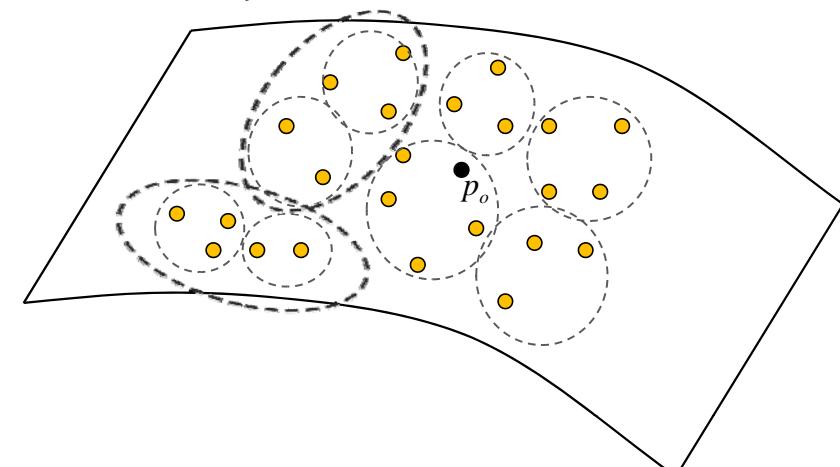
Li

$$L_o(p_o, \omega_o) \approx \frac{1}{\pi} F_t(\eta, \omega_o) F_{dt}(\eta) M_o(p_o)$$

```
if (bssrdf && octree) {
    Spectrum sigma_a = bssrdf->sigma_a();
    Spectrum sigmap_s = bssrdf->sigma_prime_s();
    Spectrum sigmap_t = sigmap_s + sigma_a;
    if (!sigmap_t.IsBlack()) {
        DiffusionReflectance Rd(sigma_a, sigmap_s,
            bssrdf->eta());
        Mo = octree->Mo(octreeBounds, p, Rd, ...);
        FresnelDielectric fresnel(1.f, bssrdf->eta());
        Ft = Spectrum(1) - fresnel.Evaluate(AbsDot(wo, n));
        float Fdt = 1.f - Fdr(bssrdf->eta());
        L += (INV_PI * Ft) * (Fdt * Mo);
    }
}
```

Mo

$$M_o(p_o) = \sum_j R_d(\|p_j - p_o\|) E_j A_j$$



SubsurfaceOctreeNode::Mo

```
Spectrum SubsurfaceOctreeNode::Mo(BBox &nodeBound,
    Point &pt, DiffusionReflectance &Rd, float maxError
{
    If extended solid angle of the node is small enough and the point is not
    inside the node, use the representative values of the node to estimate.
    float dw = sumArea / DistanceSquared(pt, p);
    if (dw < maxError && !nodeBound.Inside(pt))
        return Rd(DistanceSquared(pt, p)) * E * sumArea;

    Spectrum Mo = 0.f;
}
```



SubsurfaceOctreeNode::Mo

```
if (isLeaf)
    for (int i = 0; i < 8; ++i) {
        if (!ips[i]) break;
        Mo += Rd(DistanceSquared(pt, ips[i]->p))
            * ips[i]->E * ips[i]->area;
    }
else {
    Point pMid=.5f*nodeBound.pMin+.5f*nodeBound.pMax;
    for (int child = 0; child < 8; ++child) {
        if (!children[child]) continue;
        Bbox cBound=octreeChildBound(child,nodeBound,pMid);
        Mo+=children[child]->Mo(cBound, pt, Rd, maxError);
    }
}
return Mo;
}
```



Setting parameters



- It is remarkably unintuitive to set values of the absorption coefficient and the modified scattering coefficient.

Material	$\sigma_s' [\text{mm}^{-1}]$			$\sigma_a [\text{mm}^{-1}]$			Diffuse Reflectance			η
	R	G	B	R	G	B	R	G	B	
Apple	2.29	2.39	1.97	0.0030	0.0034	0.046	0.85	0.84	0.53	1.3
Chicken1	0.15	0.21	0.38	0.015	0.077	0.19	0.31	0.15	0.10	1.3
Chicken2	0.19	0.25	0.32	0.018	0.088	0.20	0.32	0.16	0.10	1.3
Cream	7.38	5.47	3.15	0.0002	0.0028	0.0163	0.98	0.90	0.73	1.3
Ketchup	0.18	0.07	0.03	0.061	0.97	1.45	0.16	0.01	0.00	1.3
Marble	2.19	2.62	3.00	0.0021	0.0041	0.0071	0.83	0.79	0.75	1.5
Potato	0.68	0.70	0.55	0.0024	0.0090	0.12	0.77	0.62	0.21	1.3
Skimmilk	0.70	1.22	1.90	0.0014	0.0025	0.0142	0.81	0.81	0.69	1.3
Skin1	0.74	0.88	1.01	0.032	0.17	0.48	0.44	0.22	0.13	1.3
Skin2	1.09	1.59	1.79	0.013	0.070	0.145	0.63	0.44	0.34	1.3
Spectralon	11.6	20.4	14.9	0.00	0.00	0.00	1.00	1.00	1.00	1.3
Wholemilk	2.55	3.21	3.77	0.0011	0.0024	0.014	0.91	0.88	0.76	1.3

Marble: BRDF versus BSSRDF



BRDF



BSSRDF

Marble: MCRT versus BSSRDF

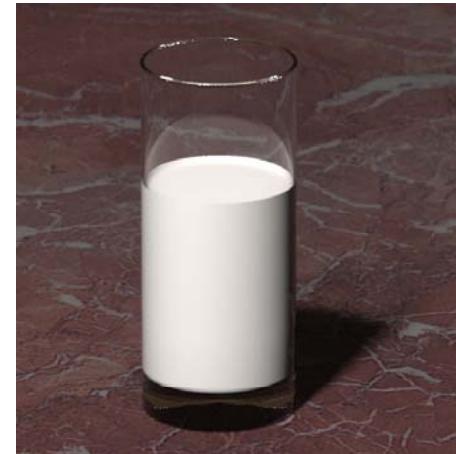


MCRT



BSSRDF

Milk



surface reflection



translucency

Skin



surface reflection



translucency